

2+1 flavor QCD with the fixed point action in the ε -regime

**Peter Hasenfratz^a, Dieter Hierl^{*b}, Vidushi Maillart^a, Ferenc Niedermayer^a,
Andreas Schäfer^b, Christof Weiermann^a, Manuel Weingart^a**

BGR (Bern-Graz-Regensburg) Collaboration

^a *Institute for Theoretical Physics, University of Bern, CH-3012 Bern, Switzerland*

^b *Institute for Theoretical Physics, University of Regensburg, D-93040 Regensburg, Germany*

E-mail: dieter.hierl@physik.uni-regensburg.de

We generated configurations with the approximate fixed-point Dirac operator D_{FP} on a 12^4 lattice with $a \approx 0.13$ fm where the scale was set by r_0 . The distributions of the low lying eigenvalues in different topological sectors were compared with those of the Random Matrix Theory which leads to a prediction of the chiral condensate.

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^{*}Speaker.

1. Introduction

The ε -expansion, where $m_\pi L \ll 1$ and $4\pi F^2 L^2 \gg 1$, is determined by Goldstone boson physics. For that reason it is an excellent tool to determine the low-energy constants of ChPT of QCD. In order to reach this regime the simulation should be done with sufficiently small quark masses. This requires a Dirac operator with good chiral behavior at the actual lattice spacing. This became possible in full QCD only recently [1, 2, 3, 4].

We consider 2 + 1 light flavor QCD applying the parametrized fixed-point (FP) action [4]. The exact FP action has nice features, most importantly, it has exact chiral symmetry. The parametrized FP action is an approximation which gave very promising results in the quenched approximation, in particular good scaling even at $a = 0.15$ fm and the spectrum of the FP Dirac operator was close to that required by chiral symmetry.

The FP Dirac operator satisfies the Ginsparg-Wilson relation:

$$D\gamma_5 + \gamma_5 D = D\gamma_5 2RD, \quad (1.1)$$

where R is a local operator living on the hypercube. The quark mass is introduced as usual for Ginsparg-Wilson operators:

$$D(m) = D + m \left(\frac{1}{2R} - \frac{1}{2} D \right). \quad (1.2)$$

The parametrized fixed-point gauge action and Dirac operator D_{FP} involve a special smearing with projection to SU(3). Therefore we cannot use a hybrid Monte Carlo algorithm. We use a partially global update procedure with nested accept/reject steps [5] relying on algorithmic developments in [6, 7, 8, 9, 10, 11, 12]. The contribution of the ~ 100 lowest lying modes to be determined is calculated exactly, the rest is treated stochastically. As a spinoff, for all the configurations in the Markov chain we have the low-lying eigenvectors which can be used in the analysis.

1.1 Remarks on the Markov chain and numerical implementation

In the algorithm the trial configurations offered in the last stochastic accept/reject step differ by a full Metropolis sweep and are accepted with $P_{\text{acc}} \approx 0.6$. We spent 13 minutes for one full accept/reject step in the Markov chain using 288 CPUs on the Altix in Munich. On this architecture it is very important to get into the cache of the single cores. Therefore we had to rewrite the whole parallelization of the existing code already running on diverse other machines while implementing it for the Altix. We use the BiCGStab(l) inverter in the stochastic estimator where l is tuned depending on the last number of inversion steps.

At the moment we have ~ 4000 configurations in the Markov chain using the partially global update algorithm [5]. In order to reduce the autocorrelation time (and the expenses of the analysis) we only take every tenth configuration minimizing the autocorrelation length and get ~ 400 pruned configurations.

1.2 Technicalities - 12^4 lattice

Our lattice spacing is $a = 0.129(5)$ fm determined from the Sommer parameter $r_0 = 0.49$ fm. Therefore our 12^4 lattice has a volume of $V \approx (1.55 \text{ fm})^4$. We calculated the additive mass shift of

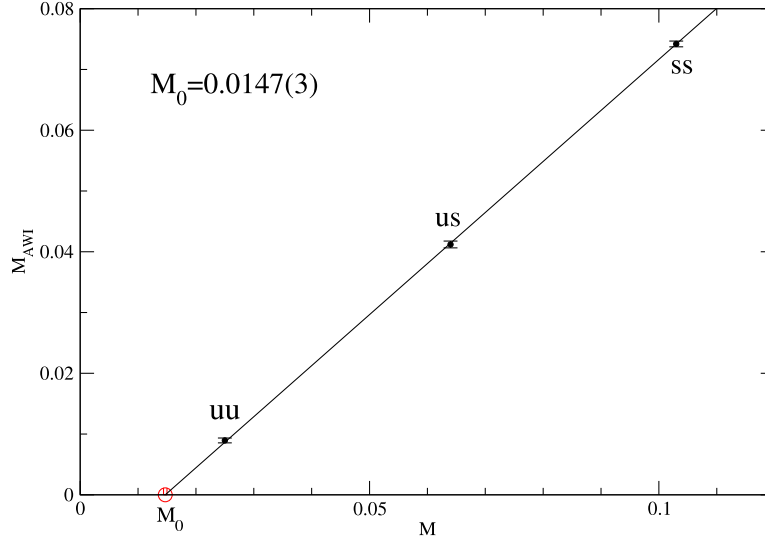


Figure 1: The AWI mass vs. the average quark mass $M = (m_f + m_{f'})/2$.

$aM_0 = 0.0147(3)$ via the axial Ward identity:

$$\partial_t \langle A_{ff'}(t) P_{ff'}(0) \rangle = 2M_{\text{AWI}} \langle P_{ff'}(t) P_{ff'}(0) \rangle, \quad \text{where } ff' = uu, us, ss. \quad (1.3)$$

In Fig. 1 we plotted the AWI mass as a function of the average quark mass $M = (m_f + m_{f'})/2$, together with a linear fit. The intercept with the horizontal axis gives the additive mass renormalization of the quark masses. One would reach the chiral limit at this value of the mass parameter in Eq. (1.2). The points uu , us and ss refer to the corresponding flavors in Eq. (1.3). Subtracting the mass shift M_0 we get for the bare masses $m_{ud} = 16$ MeV and $m_s = 137$ MeV.

2. Eigenvalues of the Dirac operator

The FP Dirac operator follows Eq. (1.1). On low-lying eigenvectors the operator $2R$ is close to 1 and we ignore it in this presentation for simplicity, i.e. we set $2R = 1$ in our discussion. Then the low-lying complex spectrum of D_{FP} lies on a circle to a good approximation, while some of the real eigenvalues are scattered away from the origin.

It is expected that the size of a effective support of the wave function and the position of the real eigenvalue are strongly correlated. Thus we introduce the inverse participation ratio (IPR) by

$$IPR = \sum_x \|\psi^{(\lambda)}(x)\|^4, \quad \text{where} \quad \|\psi^{(\lambda)}(x)\|^2 = \sum_{i=1}^{12} |\psi_i^{(\lambda)}(x)|^2 \quad (2.1)$$

for a normalized eigenvector $\psi^{(\lambda)}$. We plot the participation ratio $p = (V \times IPR)^{-1}$, where we get $p \rightarrow \mathcal{O}(1)$ for delocalized and $p \rightarrow \mathcal{O}(1/V)$ for localized eigenvectors $\psi^{(\lambda)}$.

In the analysis we used a cut $\lambda = 0.03$ (solid line in Fig. 3) in the sense that only those real eigenvalues which were below this value were considered to represent the topological charge ν . With this cut we found 328, 51 and 35 configurations for the $\nu = 0, 1, 2$ sectors, respectively.

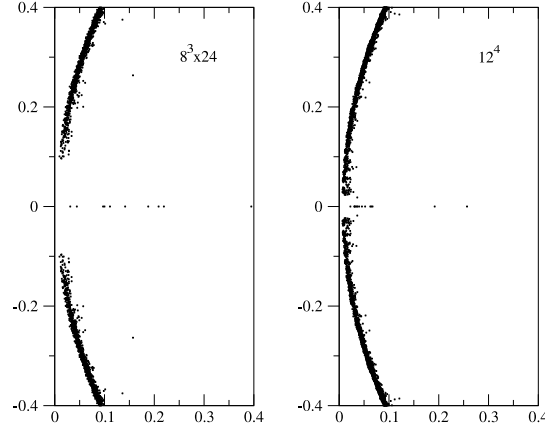


Figure 2: The left plot shows the low-lying eigenvalues of the FP Dirac operator on 50 configurations at the volume $8^3 \times 24$ and the plot on the right hand side displays the same for 12^4 configurations.

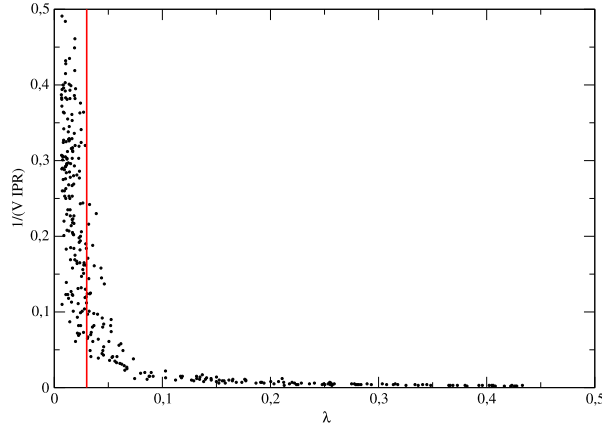


Figure 3: The relative effective volume occupied by the wave function $((V \times \text{IPR})^{-1})$ vs. the eigenvalue of D_{FP} for real modes.

This cut is sufficiently smaller than the typical gap in the complex eigenvalues (cf. Fig. 2) and corresponds to a small participation ratio (cf. Fig. 3).

3. Comparison to RMT

We use the stereographic projection of the complex eigenvalues to the imaginary axis

$$i\alpha = \frac{\lambda}{1 - \lambda/2}. \quad (3.1)$$

We denote the k -th eigenvalue in the $Q_{\text{top}} = \nu$ sector as $\alpha_{\nu k}$. Random matrix theory (RMT) predicts the distribution of the scaled eigenvalues $\xi_{\nu k}(\mu_i) = \alpha_{\nu k} \Sigma V$ depending on $\mu_i = m_i \Sigma V$, $i = 1, \dots, N_f$.

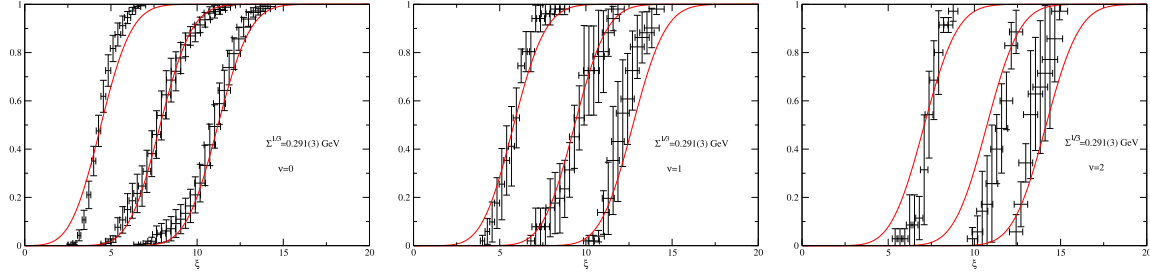


Figure 4: Cumulative distributions of $\xi_{\nu k}$ in the topological sector $\nu = 0, 1, 2$. We show the k -th smallest eigenvalues for $k = 1, 2, 3$.

The cumulative distributions of $\xi_{\nu k}$ have only one matching parameter, which is the bare condensate Σ . Fitting the distribution of the 3 lowest lying eigenvalues in the $\nu = 0, 1$ topological sector to the RMT predictions we get $\Sigma^{1/3} = 0.286(3)(9)$ GeV.

As Fig. 4 shows the distributions for different ν, k values are consistent with each other. Note, however, that the shape of the $\nu = 0, k = 1$ distribution is different from that of the RMT. The deviation in the shape could be a finite-size effect (which shows up at smallest eigenvalue, i.e. for largest wave-length), but this needs further investigations.

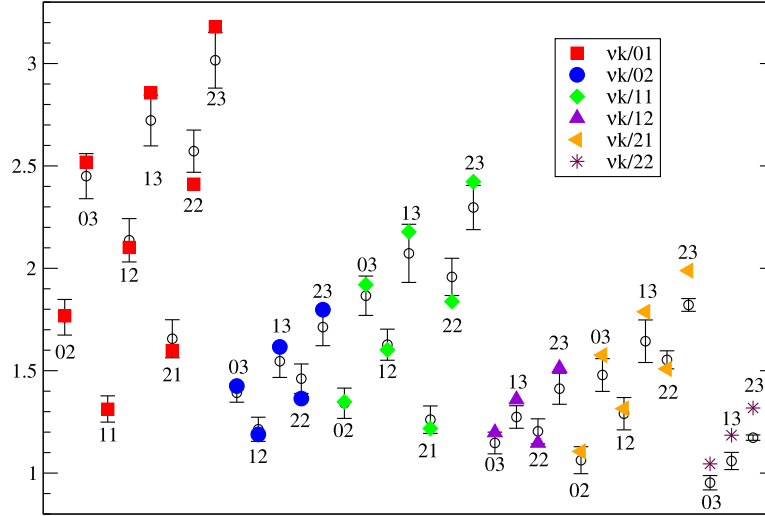


Figure 5: We calculate the ratios between different averaged eigenvalues $\langle \xi_{\nu k} \rangle$ which are denoted as νk . The different symbols denote the RMT predictions for our corresponding results.

We also compute the ratios $\langle \xi_{\nu k} \rangle / \langle \xi_{\nu' k'} \rangle$ of the mean values and compare them to the RMT predictions. We find a generally good agreement (cf. Fig. 5) even in higher topological sectors.

3.1 Finite Volume corrections

The determined Σ is the condensate of QCD in the finite volume. Thus, we have to correct the finite size effects using chiral perturbation theory with $N_f = 2$, assuming that in this correction the

s-quark contribution is negligible [4, 13]

$$\Sigma = \Sigma_\infty \left(1 + \frac{3}{2} \frac{\beta_1}{F^2 L^2} + \dots \right), \quad (3.2)$$

where $\beta_1 = 0.14046$ is a shape coefficient. After the finite size correction we get for the infinite-volume bare condensate

$$\Sigma_\infty^{1/3} = 0.255(3)(9) \text{ GeV}. \quad (3.3)$$

4. Summary and Outlook

Simulating the parametrized FP action with partially global update procedure we were able to reach sufficiently small quark masses to study the ε -regime in the 2+1 flavor QCD. The distribution of the low-lying eigenvalues in different topological sectors is in good agreement with the RMT predictions, although the identification of the topological charge is somewhat ambiguous.

The next steps are to compute the low energy constant F from the $\langle PP \rangle$ correlator, to calculate the corresponding Z factors and to complete a lattice simulation on a $12^3 \times 24$ lattice with comparable lattice spacing in the δ -regime.

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